# Light meson spectroscopy with the KLOE experiment.

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#### Abstract

In this paper we describe the status of the analyses in progress on light meson spectroscopy in the KLOE experiment. We present the analyses of  $\phi$  decays into  $f_0(980)\gamma$  and  $a_0(980)\gamma$ , the Dalitz plot analysis of the  $\eta \to \pi^+\pi^-\pi^0$  decay, the branching ratio measurement of  $\eta \to \pi^0\gamma\gamma$ , the upper limits on  $Br(\eta \to 3\gamma)$  and  $Br(\eta \to \pi^+\pi^-)$ , the measurement of the ratio  $Br(\phi \to \eta'\gamma)/Br(\phi \to \eta\gamma)$  and  $\phi$  leptonic width measurements.

### 1 Introduction

The KLOE detector [1], operates at the Frascati  $e^+e^-$  collider DA $\Phi$ NE [2], which runs at a CM energy W equal to the  $\phi$ -meson mass,  $W{\sim}1019.5$  MeV. The analyses presented here are based on data collected in the years 2001 and 2002 for an integrated luminosity of  $\sim 450~pb^{-1}$  corresponding to 1.5 billions of  $\phi$  and 20 millions of  $\eta$  mesons [Br( $\phi \rightarrow \eta \gamma$ )  $\sim 1.3\%$  [3]]. This means that KLOE can also study  $\eta$  physics in a clean environment with high statistic.

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## 2 Search for $\phi \rightarrow f_0 \gamma$ in $\pi^+ \pi^- \gamma$ events.

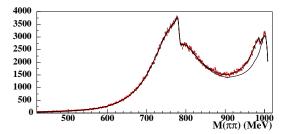
The  $\phi$  radiative decays to scalar mesons,  $\phi \rightarrow S\gamma$ , give significant insight in the assessment of the nature of lower mass scalar mesons [4]. With the KLOE experiment the decays  $\phi \to f_0(980)\gamma$  and  $\phi \to a_0(980)\gamma$  are searched for in  $\pi^0\pi^0\gamma$  and  $\eta\pi^0\gamma$  [5,6] final states respectively. Moreover the  $f_0(980)$  signal is also searched for  $\pi^+\pi^-\gamma$  events with a photon at large angle. The search for this signal is characterized by the presence of a huge irreducible background due to the initial state radiation (ISR), to  $e^+e^- \to \pi^+\pi^-\gamma$  (FSR) and  $\phi \rightarrow \rho^{\pm}(\rightarrow \pi^{\pm}\gamma)\pi^{\pm}$ . The f<sub>0</sub> events are searched for in the large photon angle region  $45^{o} < \theta < 135^{o}$  to reduce ISR background. The f<sub>0</sub> signal appears as a bump in the  $\pi^+\pi^-$  invariant mass  $M_{\pi\pi}$  spectrum around 980 MeV. Fig.1 (top) shows the spectrum obtained at  $\sqrt{s} = M_{\phi}$ .

An overall fit to the spectrum has been done by applying the following formula:

$$\frac{dN}{dM_{\pi\pi}} = \left[ \left( \frac{d\sigma}{dM_{\pi\pi}} \right)_{ISR} + \left( \frac{d\sigma}{dM_{\pi\pi}} \right)_{FSR+f_0} + \left( \frac{d\sigma}{dM_{\pi\pi}} \right)_{\rho\pi} \right] \times L \times \epsilon(M_{\pi\pi})$$

with L the integrated luminosity and  $\epsilon(M_{\pi\pi})$  the selection efficiency as a function of  $M_{\pi\pi}$ . The  $f_0$  amplitude is taken from the kaon-loop approach [7]. A forward-backward asymmetry  $A = \frac{N^+(\theta > 90^\circ) - N^+(\theta < 90^\circ)}{N^+(\theta > 90^\circ) + N^+(\theta < 90^\circ)}$  is expected, due to the interference between FSR

and ISR[8]. Fig.1 (bottom) shows the asymmetry as a function of  $M_{\pi\pi}$  compared to the prediction based on the ISR-FSR interference alone. A significant deviation from the prediction is observed in the high mass region clearly due to the  $f_0$  contribution.



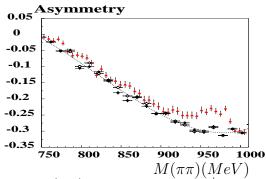


Fig. 1. (top)  $M_{\pi\pi}$  spectrum of  $\pi^+\pi^-\gamma$ . The upper (lower) curves are the result of the fit and the estimated background due to ISR, FSR and  $\rho\pi$ . (bottom) Forward-Backward asymmetry A as a function of  $M_{\pi\pi}$ . The curve and the black points are the Montecarlo expectations based on the interference between FSR and ISR only. The experimental data are reported as triangles.

## 3 Dynamics of $\eta \to \pi^+ \pi^- \pi^0$

The dynamics of the  $\eta \to \pi^+ \pi^- \pi^0$  decay has been studied with a Dalitz plot analysis. The conventional variables X and Y are defined as:  $X = \sqrt{3} \frac{T_+ - T_-}{Q_\eta}, Y = \frac{3T_0}{Q_\eta} - 1$ , where  $Q_\eta = m_\eta - 2m_{\pi^+} - m_{\pi^0}$  and  $T_+$ ,  $T_-$  and  $T_0$  are the kinetic energies of the particles. The measured distribu-

tion has been fitted as:  $|A(X,Y)|^2 \simeq (1+aY+bY^2+cX+dX^2+eXY+...)$ . C-parity conservation prevents odd powers in X being present in the expansion: thus parameters c and e should be zero as confirmed by our fit. The results of the fit are shown in table 3 Efficiency is  $\sim 36$  % over

Table 1 Fitted parameters  $P(\chi^2) = 52 \%$  of  $\eta \to 3\pi$  Dalitz plot.

a	b	c
$-1.075 \pm .008$	$.118 \pm .009$	$0005 \pm .004$
d	e	f
$.049 \pm .008$	$004 \pm .01$	$.13 \pm .02$

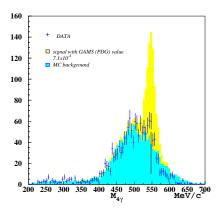
the whole Dalitz plot. The evaluation of systematic effects is under completion.

# $\begin{array}{cccc} \textbf{4} & \textbf{Rare and forbidden } \eta \textbf{ decays} \\ & (\eta \ \rightarrow \ \pi^0 \gamma \gamma, \! \eta \ \rightarrow \ \pi^+ \pi^-, \! \eta \ \rightarrow \\ & \gamma \gamma \gamma) \end{array}$

The  $\eta \to \pi^0 \gamma \gamma$  decay is interesting to test the Chiral Perturbation Theory prediction for the branching ratio and  $m_{\gamma\gamma}$  spectrum[9,10]. The most accurate measurement for the branching ratio[11] is, infact, far from any theoretical prediction for this decay based on ChPT. Recently a new measurement has been performed [12] giving a much lower value than the previous one, with a larger error. All previuos experiments were done at hadron machines, using mainly  $\pi^- p \rightarrow \eta n$ , and are largely dominated by  $\pi^0\pi^0$  background and geometrical acceptance. KLOE performs a measurement of competitive precision in a cleaner environment.

Furthermore, it has different background topologies and experimental systematics. The signal is searched for by looking for a  $\pi^0 \gamma \gamma \gamma$  topology, where the further  $\gamma$  comes from  $\phi \to \eta \gamma$ . Five prompt clusters are required and an overall kinematic fit requiring  $\pi^0$  mass is performed. The clusters energy must be > 30 MeVand azimutal angle  $> 21^{\circ}$  to reject fake clusters coming from machine background. The dominant background channel is  $\eta \to 3\pi^0$  that had been reduced with several topological cut. With this selection we obtain an efficiency of 5.7%. To give an idea of the sensitivity, In fig. 2 we compare the  $M(4\gamma)$  data spectrum with MC based predictions of signal and background in two hypothesys for the size of signal: one based on PDG value[11] and one based on the recent CB result[12]. It is evident that our data are incompatible with [11] and are marginally in agreement with [12]. The background simulation and the efficiency for the signal is still under study.

 $\eta \to 3\gamma$  decay is C violating. It is a sensitive test of C violation in the strong and electromagnetic interactions. For the details of this analysis see ref.[13]. The KLOE result for the branching ratio is:  $Br(\eta \to \gamma\gamma\gamma) \le 1.6 \times 10^{-5}$  @90% C.L. This limit is the best experimental limit for this decay. The expected branching ratio from the Standard Model is  $\le 10^{-12}$  [14], so any discovery of a larger decay rate would be a clear signal of Standard Model deviation.



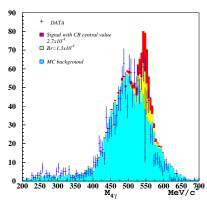


Fig. 2.  $M(4\gamma)$ , the spectra expected from the GAMS[11] and Crystall Ball[12] measurement are shown. In the lower plot we show also the expected spectrum for a  $Br \sim 1/2$  of C.B. result.

 $\eta \to \pi^+\pi^-$  decay is P and CP violating. This decay is allowed as a weak direct CP violating decay with a very low branching ratio:  $BR(\eta \to \pi^+\pi^-) \sim 10^{-27}$  [15]. Therefore the detection of this decay at an accessible level would be a signal of P and CP violation not explainable in the Standard Model framework. The latest published [16] direct search of this decay has given the following 90% C.L. upper limit:  $BR(\eta \to \pi^+\pi^-) < 3.3 \times 10^{-4}$ . In KLOE the signal is searched for the  $M(\eta)$  region of the  $\pi^+\pi^-$  invariant mass spectrum of  $\pi^+\pi^-\gamma$ events selected according to the  $f_0(980) \rightarrow \pi^+\pi^-$  analysis described before (see fig. 1). The signal efficiency is:  $\epsilon_s=16.6\%$ . The expected signal has a Gaussian shape with a mass resolution of 1.33 MeV. No signal is observed. The background is determined by fitting the theoretical model for  $\pi^+\pi^-\gamma$  sample to the full spectrum. In order to determine an upper limit, we have added to this background a Gaussian function representing the signal multiplied by a constant  $N_s$ . We obtain:  $N_s = -8 \pm 24$ . The 90% confidence level upper limit on the number of events is obtained using the tables in [17]:  $N_s$  < 33. The branching ratio is  $BR(\eta \rightarrow \pi^+\pi^-) = \frac{N_s}{\epsilon_s N_\eta}$ with  $N_{\eta}$  the number of  $\eta$  in the sample  $(1.55 \times 10^7)$ . The 90% C.L. upper limit is:  $BR(\eta \to \pi^+\pi^-)$  <  $1.3 \times 10^{-5}$ . It improves by a factor  $\sim$ 30 the current PDG limit.

### 5 $\eta$ - $\eta'$ mixing

Here we present the  $R = \frac{\Gamma(\phi \to \eta' \gamma)}{\Gamma(\phi \to \eta \gamma)}$ measurement. The  $\eta'$  is identified via the decays:  $\phi \to \eta' \gamma$ ;  $\eta' \to \pi^+ \pi^- \eta$ ;  $\eta \to \pi^0 \pi^0 \pi^0$  and the decays  $\phi \to$  $\eta'\gamma~;~\eta'~\to\pi^0\,\pi^0\,\eta~;~\eta~\to\pi^+\pi^-\pi^0~.$ The final state is thus charachterized by two charged pions and seven photons, and has no physics background with the same topology in KLOE. After background subtraction we observe  $3405 \pm 61 \pm 31 \phi \rightarrow \eta' \gamma$  events. We normalize to the number of observed  $\eta \to \pi^0 \pi^0 \pi^0$  decays in the same runs to obtain a preliminary measurement of the ratio of BR's:  $R = (4.9 \pm 0.1_{stat} \pm 0.2_{syst}) \times 10^{-3}$ .

This result compares favourably with our previous estimate [18] (which already dominates the world average [3]) but with considerably improved accuracy.

## 6 A new measurement of the $\phi$ leptonic width.

KLOE has performed a new measurement of the  $\phi$  leptonic widths  $\Gamma_{ll}$  with  $l = e, \mu$  [19], using the two data samples taken below ( $\sqrt{s}=1017 \text{ MeV}$ ) and above  $(\sqrt{s}=1022 \text{ MeV})$  the  $\phi$ peak together with the data taken at the  $\phi$  peak. The dependences on  $\sqrt{s}$  of the forward-backward asymmetry of Bhabha events  $A_{FB}$  and of the  $e^+e^- \rightarrow \mu^+\mu^-$  cross-section  $\sigma(\mu\mu)$  around the  $\phi$  peak are sensitive to the value of  $\Gamma_{ee}$  and  $\sqrt{\Gamma_{ee}\Gamma_{\mu\mu}}$ respectively. We measure the  $\phi$  mass  $M_{\phi}$ , the forward-backward asymmetry at  $W = M_{\phi} A_{FB}^{0}$ , and finally  $\Gamma_{ee}$ . The result for  $\Gamma_{ee}$  is:  $\Gamma_{ee} =$  $1.32 \pm 0.05_{stat} \pm 0.03_{syst}$  keV The result for  $\sqrt{\Gamma_{ee}\Gamma_{\mu\mu}}$  is :  $\sqrt{\Gamma_{ee}\Gamma_{\mu\mu}} = 1.320 \pm 0.018_{stat} \pm 0.017_{syst}$  keV. The two results are in good agreement consistently with lepton universality. Combining them we get:  $\Gamma_{ll} = 1.320 \pm 0.017_{stat} \pm 0.015_{syst} \, keV$ with a total uncertainty below 2 %. We point out that the value of  $\Gamma_{ee}$ is necessary for  $\phi$  decay branching ratio measurements, and play also a role in the evaluation of the hadronic contribution to vacuum polarization [20].

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